

# Introduction to Quantum Information Processing

QIC 710 / CS 768 / PH 767 / CO 681 / AM 871

## Lectures 10–11 (2019)

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# More state distinguishing problems

# More state distinguishing problems

Which of these states are distinguishable? Divide them into equivalence classes:

1.  $|0\rangle + |1\rangle$

2.  $-|0\rangle - |1\rangle$

3.  $\begin{cases} |0\rangle & \text{with prob. } \frac{1}{2} \\ |1\rangle & \text{with prob. } \frac{1}{2} \end{cases}$

4.  $\begin{cases} |0\rangle + |1\rangle & \text{with prob. } \frac{1}{2} \\ |0\rangle - |1\rangle & \text{with prob. } \frac{1}{2} \end{cases}$

5.  $\begin{cases} |0\rangle & \text{with prob. } \frac{1}{2} \\ |0\rangle + |1\rangle & \text{with prob. } \frac{1}{2} \end{cases}$

6.  $\begin{cases} |0\rangle & \text{with prob. } \frac{1}{4} \\ |1\rangle & \text{with prob. } \frac{1}{4} \\ |0\rangle + |1\rangle & \text{with prob. } \frac{1}{4} \\ |0\rangle - |1\rangle & \text{with prob. } \frac{1}{4} \end{cases}$

7. The first qubit of  $|01\rangle - |10\rangle$

Answers later on ...

**This is a probabilistic mixed state**

# Density matrix formalism

# Density matrices (1)

Until now, we've represented quantum states as **vectors** (e.g.  $|\psi\rangle$ , and all such states are called **pure states**)

An alternative way of representing quantum states is in terms of **density matrices** (a.k.a. **density operators**)

The density matrix of a pure state  $|\psi\rangle$  is the matrix  $\rho = |\psi\rangle\langle\psi|$

**Example:** the density matrix of  $\alpha|0\rangle + \beta|1\rangle$  is

$$\rho = \begin{bmatrix} \alpha \\ \beta \end{bmatrix} \begin{bmatrix} \alpha^* & \beta^* \end{bmatrix} = \begin{bmatrix} |\alpha|^2 & \alpha\beta^* \\ \alpha^*\beta & |\beta|^2 \end{bmatrix}$$

# Density matrices (2)

How do quantum operations work using density matrices?

**Effect of a unitary operation on a density matrix:**

applying  $U$  to  $\rho$  yields  $U\rho U^\dagger$

(this is because the modified state is  $U|\psi\rangle\langle\psi|U^\dagger$ )

**Effect of a measurement on a density matrix:**

measuring state  $\rho$  with respect to the basis  $|\varphi_1\rangle, |\varphi_2\rangle, \dots, |\varphi_d\rangle$ , yields the  $k^{\text{th}}$  outcome with probability  $\langle\varphi_k|\rho|\varphi_k\rangle$

(this is because  $\langle\varphi_k|\rho|\varphi_k\rangle = \langle\varphi_k|\psi\rangle\langle\psi|\varphi_k\rangle = |\langle\varphi_k|\psi\rangle|^2$ )

—and the state collapses to  $|\varphi_k\rangle\langle\varphi_k|$

# Density matrices (3)

A probability distribution on pure states is called a ***mixed state***:

$$\left( (|\psi_1\rangle, p_1), (|\psi_2\rangle, p_2), \dots, (|\psi_d\rangle, p_d) \right)$$

The ***density matrix*** associated with such a mixed state is:

$$\rho = \sum_{k=1}^d p_k |\psi_k\rangle\langle\psi_k|$$

**Example:** the density matrix for  $\left( (|0\rangle, \frac{1}{2}), (|1\rangle, \frac{1}{2}) \right)$  is:

$$\frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} + \frac{1}{2} \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

**Question:** what is the density matrix of

$$\left( (|0\rangle + |1\rangle, \frac{1}{2}), (|0\rangle - |1\rangle, \frac{1}{2}) \right) ?$$

# Density matrices (4)

How do quantum operations work for these *mixed* states?

**Effect of a unitary operation on a density matrix:**

applying  $U$  to  $\rho$  *still* yields  $U\rho U^\dagger$

This is because the modified state is:

$$\sum_{k=1}^d p_k U |\psi_k\rangle \langle \psi_k| U^\dagger = U \left( \sum_{k=1}^d p_k |\psi_k\rangle \langle \psi_k| \right) U^\dagger = U \rho U^\dagger$$

**Effect of a measurement on a density matrix:**

measuring state  $\rho$  with respect to the basis  $|\varphi_1\rangle, |\varphi_2\rangle, \dots, |\varphi_d\rangle$ , *still* yields the  $k^{\text{th}}$  outcome with probability  $\langle \varphi_k | \rho | \varphi_k \rangle$

**Why?**



# Recap: density matrices

## Quantum operations in terms of density matrices:

- Applying  $U$  to  $\rho$  yields  $U\rho U^\dagger$
- Measuring state  $\rho$  with respect to the basis  $|\varphi_1\rangle, |\varphi_2\rangle, \dots, |\varphi_d\rangle$ , yields:  $k^{\text{th}}$  outcome with probability  $\langle \varphi_k | \rho | \varphi_k \rangle$   
—and causes the state to collapse to  $|\varphi_k\rangle\langle \varphi_k|$

Since these are expressible in terms of density matrices alone (independent of any specific probabilistic mixtures), states with identical density matrices are ***operationally indistinguishable***

Return to state distinguishing  
problems ...

# State distinguishing problems (1)

The **density matrix** of the mixed state

$((|\psi_1\rangle, p_1), (|\psi_2\rangle, p_2), \dots, (|\psi_d\rangle, p_d))$  is: 
$$\sum_{k=1}^d p_k |\psi_k\rangle \langle \psi_k|$$

**Examples (from earlier in lecture):**

1. & 2.  $|0\rangle + |1\rangle$  and  $-|0\rangle - |1\rangle$  both have 
$$\rho = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}$$

3.  $\left\{ \begin{array}{l} |0\rangle \text{ with prob. } \frac{1}{2} \\ |1\rangle \text{ with prob. } \frac{1}{2} \end{array} \right.$

4.  $\left\{ \begin{array}{l} |0\rangle + |1\rangle \text{ with prob. } \frac{1}{2} \\ |0\rangle - |1\rangle \text{ with prob. } \frac{1}{2} \end{array} \right.$

6.  $\left\{ \begin{array}{l} |0\rangle \text{ with prob. } \frac{1}{4} \\ |1\rangle \text{ with prob. } \frac{1}{4} \\ |0\rangle + |1\rangle \text{ with prob. } \frac{1}{4} \\ |0\rangle - |1\rangle \text{ with prob. } \frac{1}{4} \end{array} \right.$

$$\rho = \frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

# State distinguishing problems (2)

## Examples (continued):

5.  $\begin{cases} |0\rangle & \text{with prob. } \frac{1}{2} \\ |0\rangle + |1\rangle & \text{with prob. } \frac{1}{2} \end{cases}$

$$\text{has: } \rho = \frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} + \frac{1}{2} \begin{bmatrix} 1/2 & 1/2 \\ 1/2 & 1/2 \end{bmatrix} = \begin{bmatrix} 3/4 & 1/2 \\ 1/2 & 1/4 \end{bmatrix}$$

7. The first qubit of  $|01\rangle - |10\rangle$  ...? (later)

# Characterizing density matrices

Three properties of  $\rho$  :

- $\text{Tr}\rho = 1$  ( $\text{Tr}M = M_{11} + M_{22} + \dots + M_{dd}$ )
- $\rho = \rho^\dagger$  (i.e.  $\rho$  is *Hermitian*)
- $\langle \varphi | \rho | \varphi \rangle \geq 0$ , for all states  $|\varphi\rangle$  (i.e.  $\rho$  is *positive semidefinite*)

$$\rho = \sum_{k=1}^d p_k |\psi_k\rangle\langle\psi_k|$$

Moreover, for **any** matrix  $\rho$  satisfying the above properties, there exists a probabilistic mixture whose density matrix is  $\rho$

**Exercise:** show this

# Taxonomy of various normal matrices

# Normal matrices

**Definition:** A matrix  $M$  is *normal* if  $M^\dagger M = M M^\dagger$

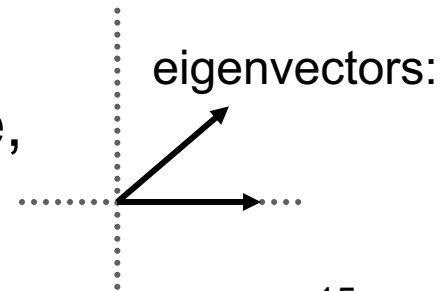
**Theorem:**  $M$  is normal iff there exists a unitary  $U$  such that  $M = U^\dagger D U$ , where  $D$  is diagonal (i.e. unitarily diagonalizable)

$$D = \begin{bmatrix} \lambda_1 & 0 & \cdots & 0 \\ 0 & \lambda_2 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & \lambda_d \end{bmatrix}$$

Examples of *ab*normal matrices:

$\begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix}$  is not even diagonalizable

$\begin{bmatrix} 1 & 1 \\ 0 & 2 \end{bmatrix}$  is diagonalizable, but not unitarily



# Unitary and Hermitian matrices

**Normal:**  $M = \begin{bmatrix} \lambda_1 & 0 & \cdots & 0 \\ 0 & \lambda_2 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & \lambda_d \end{bmatrix}$  with respect to **some** orthonormal basis

**Unitary:**  $M^\dagger M = I$  which implies  $|\lambda_k|^2 = 1$ , for all  $k$

**Hermitian:**  $M = M^\dagger$  which implies  $\lambda_k \in \mathbb{R}$  for all  $k$

**Question:** which matrices are both unitary **and** Hermitian?

**Answer:** reflections ( $\lambda_k \in \{+1, -1\}$ , for all  $k$ )



# Positive semidefinite

**Positive semidefinite:** Hermitian and  $\lambda_k \geq 0$ , for all  $k$

**Theorem:**  $M$  is positive semidefinite iff  $M$  is Hermitian and, for all  $|\varphi\rangle$ ,  $\langle \varphi | M | \varphi \rangle \geq 0$

**(Positive definite:**  $\lambda_k > 0$ , for all  $k$ )

# Projectors and density matrices

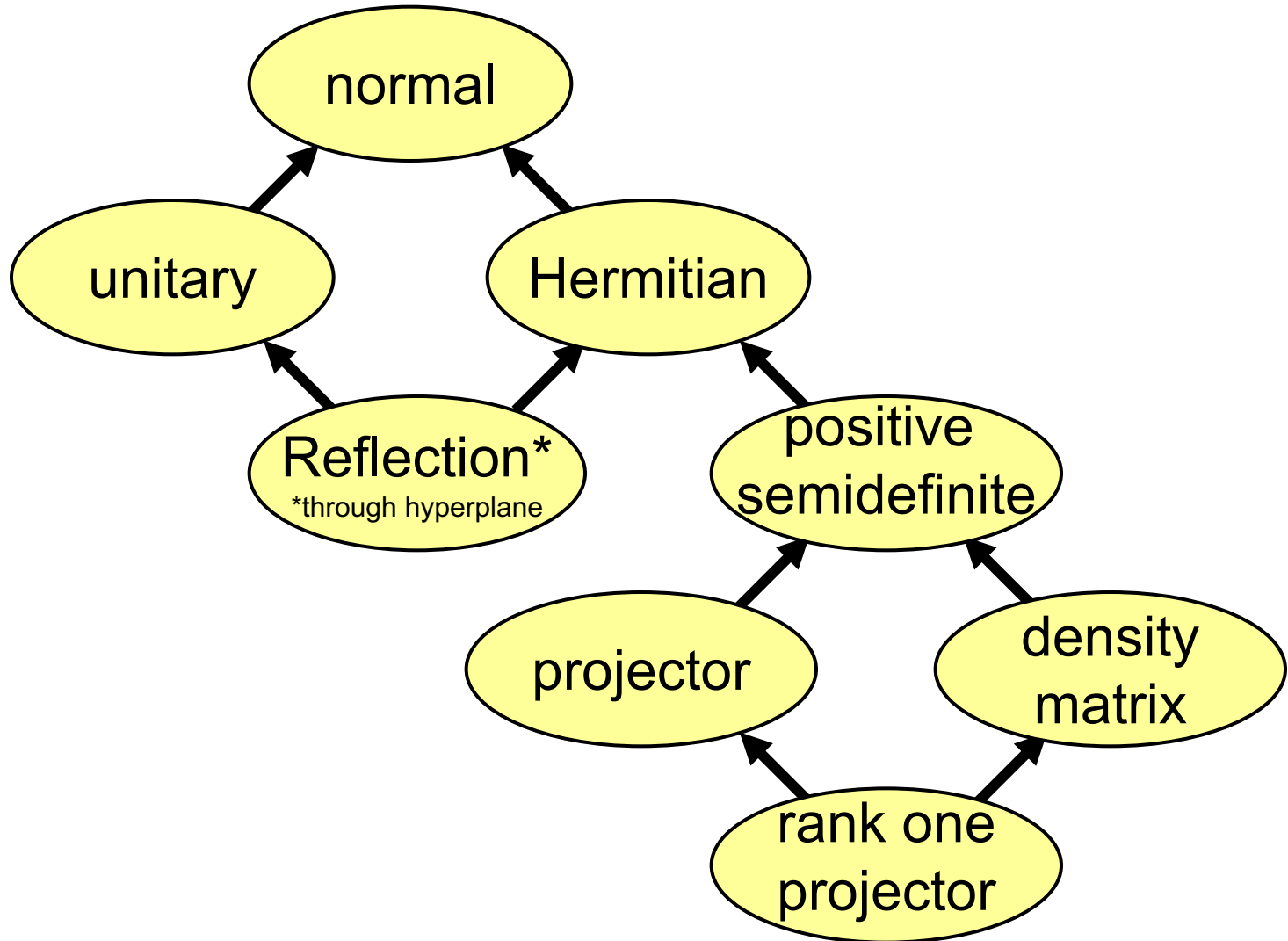
**Projector:** Hermitian and  $M^2 = M$ , which implies that  $M$  is positive semidefinite and  $\lambda_k \in \{0,1\}$ , for all  $k$

**Density matrix:** positive semidefinite and  $\text{Tr } M = 1$ , so  $\sum_{k=1}^d \lambda_k = 1$

**Question:** which matrices are both projectors *and* density matrices?

**Answer:** rank-1 projectors ( $\lambda_k = 1$  if  $k = j$ ; otherwise  $\lambda_k = 0$ )

# Taxonomy of normal matrices



# Bloch sphere for qubits

# Bloch sphere for qubits (1)

Consider the set of all 2x2 density matrices  $\rho$

They have a nice representation in terms of the **Pauli matrices**:

$$\sigma_x = X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \quad \sigma_z = Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \quad \sigma_y = Y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$$

Note that these matrices—combined with  $I$ —form a **basis** for the vector space of all 2x2 matrices

We will express density matrices  $\rho$  in this basis

**Note:** coefficient of  $I$  must be  $\frac{1}{2}$ , since  $X, Y, Z$  are traceless

# Bloch sphere for qubits (2)

We will express  $\rho = \frac{I + c_x X + c_y Y + c_z Z}{2}$

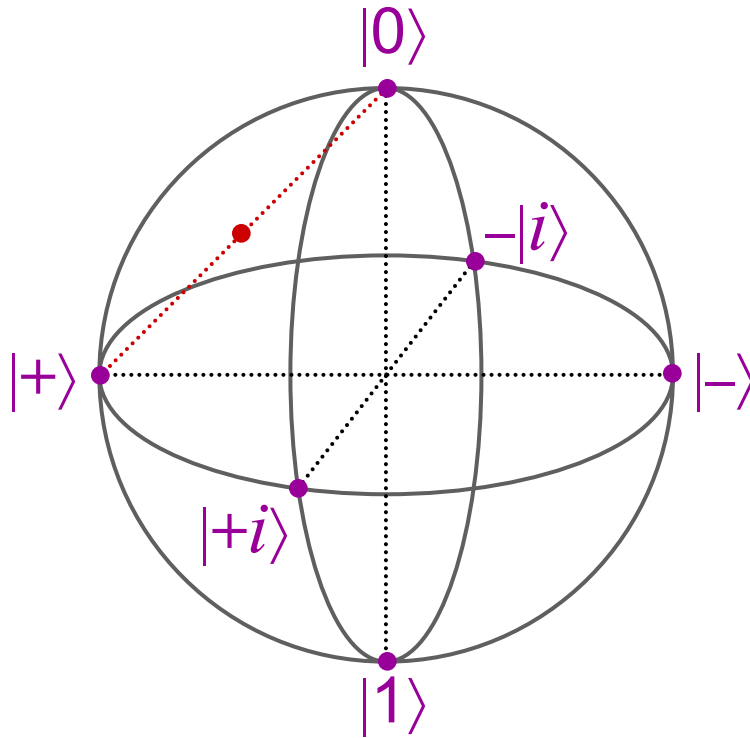
First consider the case of pure states  $|\psi\rangle\langle\psi|$ , where, without loss of generality,  $|\psi\rangle = \cos(\theta)|0\rangle + e^{2i\phi}\sin(\theta)|1\rangle$  ( $\theta, \phi \in [0, \pi]$ )

$$\rho = \begin{bmatrix} \cos^2\theta & e^{-i2\phi}\cos\theta\sin\theta \\ e^{i2\phi}\cos\theta\sin\theta & \sin^2\theta \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1 + \cos(2\theta) & e^{-i2\phi}\sin(2\theta) \\ e^{i2\phi}\sin(2\theta) & 1 - \cos(2\theta) \end{bmatrix}$$

Therefore  $c_z = \cos(2\theta)$ ,  $c_x = \cos(2\phi)\sin(2\theta)$ ,  $c_y = \sin(2\phi)\sin(2\theta)$

These are **polar coordinates** of a unit vector  $(c_x, c_y, c_z) \in \mathbb{R}^3$

# Bloch sphere for qubits (3)



$$|+\rangle = |0\rangle + |1\rangle$$

$$|-\rangle = |0\rangle - |1\rangle$$

$$|+i\rangle = |0\rangle + i|1\rangle$$

$$|-i\rangle = |0\rangle - i|1\rangle$$

Note that *orthogonal* corresponds to *antipodal* here

Pure states are on the surface, and mixed states are inside (being weighted averages of pure states)

# Distinguishing mixed states



# Distinguishing mixed states (1)

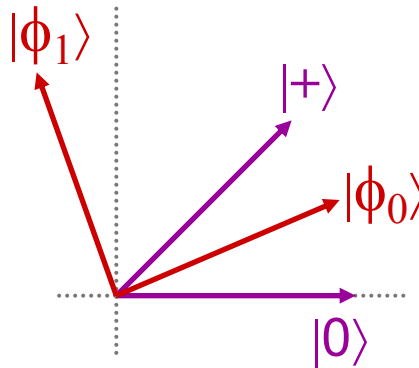
What's the best distinguishing strategy between these two mixed states?

$$\begin{cases} |0\rangle & \text{with prob. } 1/2 \\ |0\rangle + |1\rangle & \text{with prob. } 1/2 \end{cases}$$

$$\rho_1 = \begin{bmatrix} 3/4 & 1/2 \\ 1/2 & 1/4 \end{bmatrix}$$

$\rho_1$  also arises from this orthogonal mixture:

$$\begin{cases} |\phi_0\rangle & \text{with prob. } \cos^2(\pi/8) \\ |\phi_1\rangle & \text{with prob. } \sin^2(\pi/8) \end{cases}$$



$$\begin{cases} |0\rangle & \text{with prob. } 1/2 \\ |1\rangle & \text{with prob. } 1/2 \end{cases}$$

$$\rho_2 = \frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

... as does  $\rho_2$  from:

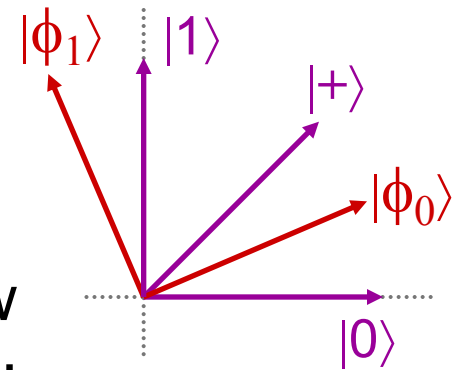
$$\begin{cases} |\phi_0\rangle & \text{with prob. } 1/2 \\ |\phi_1\rangle & \text{with prob. } 1/2 \end{cases}$$

# Distinguishing mixed states (2)

We've effectively found an orthonormal basis  $|\phi_0\rangle, |\phi_1\rangle$  in which both density matrices are diagonal:

$$\rho'_2 = \begin{bmatrix} \cos^2(\pi/8) & 0 \\ 0 & \sin^2(\pi/8) \end{bmatrix} \quad \rho'_1 = \frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

Rotating  $|\phi_0\rangle, |\phi_1\rangle$  to  $|0\rangle, |1\rangle$  the scenario can now be examined using classical probability theory:



Distinguish between two **classical** coins, whose probabilities of “heads” are  $\cos^2(\pi/8)$  and  $1/2$  respectively (details: exercise)

**Question:** what do we do if we aren't so lucky to get two density matrices that are simultaneously diagonalizable?

**general quantum operations**

more commonly known as

**quantum channels**

# General quantum operations (1)

Also known as:

“quantum channels”

“completely positive trace preserving maps”,

“admissible operations”

Let  $A_1, A_2, \dots, A_m$  be matrices satisfying  $\sum_{j=1}^m A_j^\dagger A_j = I$

Then the mapping  $\rho \mapsto \sum_{j=1}^m A_j \rho A_j^\dagger$  is a general quantum op

**Note:**  $A_1, A_2, \dots, A_m$  do not have to be square matrices

**Example 1 (unitary op):** applying  $U$  to  $\rho$  yields  $U\rho U^\dagger$

# General quantum operations (2)

**Example 2 (decoherence):** let  $A_0 = |0\rangle\langle 0|$  and  $A_1 = |1\rangle\langle 1|$

This quantum op maps  $\rho$  to  $|0\rangle\langle 0|\rho|0\rangle\langle 0| + |1\rangle\langle 1|\rho|1\rangle\langle 1|$

For  $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$ ,

$$\begin{bmatrix} |\alpha|^2 & \alpha\beta^* \\ \alpha^*\beta & |\beta|^2 \end{bmatrix} \mapsto \begin{bmatrix} |\alpha|^2 & 0 \\ 0 & |\beta|^2 \end{bmatrix}$$

Corresponds to measuring  $\rho$  “without looking at the outcome”

After looking at the outcome,  $\rho$  becomes  $\begin{cases} |0\rangle\langle 0| & \text{with prob. } |\alpha|^2 \\ |1\rangle\langle 1| & \text{with prob. } |\beta|^2 \end{cases}$

# General quantum operations (3)

## Example 3

$$\text{Let } A_0 = I \otimes \langle 0 | = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{bmatrix} \text{ and } A_1 = I \otimes \langle 1 | = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

- Any state of the form  $\rho \otimes \sigma$  (product state) becomes  $\rho$
- State  $\left(\frac{1}{\sqrt{2}}|00\rangle + \frac{1}{\sqrt{2}}|11\rangle\right)\left(\frac{1}{\sqrt{2}}\langle 00| + \frac{1}{\sqrt{2}}\langle 11|\right)$  becomes  $\frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$

It's the same density matrix as for  $\left(\frac{1}{2}, |0\rangle\right), \left(\frac{1}{2}, |1\rangle\right)$

- Corresponds to “discarding the second register”

The operation is called the **partial trace**  $\text{Tr}_2 \rho$